

## Key Concepts for this section

- 1: Lorentz force law, Field, Maxwell's equation
- 2: Ion Transport, Nernst-Planck equation
- 3: (Quasi)electrostatics, potential function,
- 4: Laplace's equation, Uniqueness
- 5: Debye layer, electroneutrality

### **Goals of Part II:**

- (1) Understand when and why electromagnetic (E and B) interaction is relevant (or not relevant) in biological systems.
- (2) Be able to analyze quasistatic electric fields in 2D and 3D.

$$\vec{E} = -\nabla\Phi \quad \nabla \cdot (\epsilon\vec{E}) = \nabla \cdot (-\epsilon\nabla\Phi) = \rho_e$$

$$\nabla^2\Phi = -\frac{\rho_e}{\epsilon} \quad (\text{Poisson's Equation})$$

However, biomolecules in the system do not generate E-field, since they are shielded by counterions (electroneutrality).....

It all comes down to solving.....  $\nabla^2\Phi = 0$  (*Laplace's Equation*)

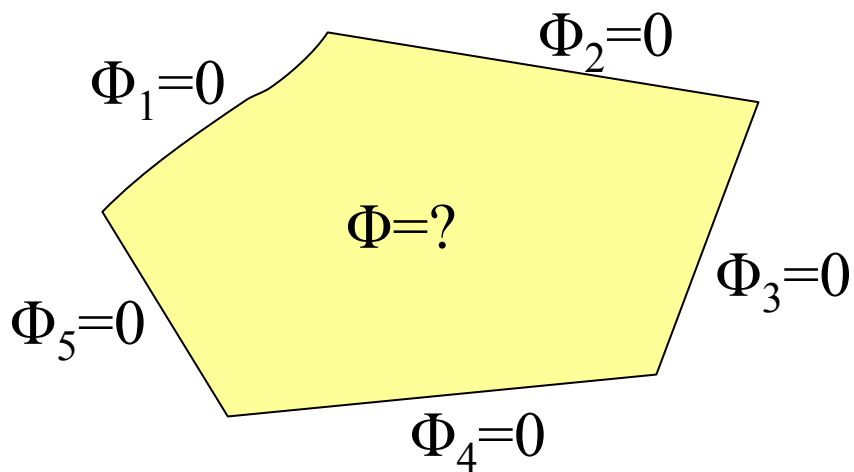
$$\nabla^2 c = \frac{\partial c}{\partial t} \quad (\text{Fick's second law}) \quad \longrightarrow \quad \nabla^2 c = 0$$

(steady-state diffusion)

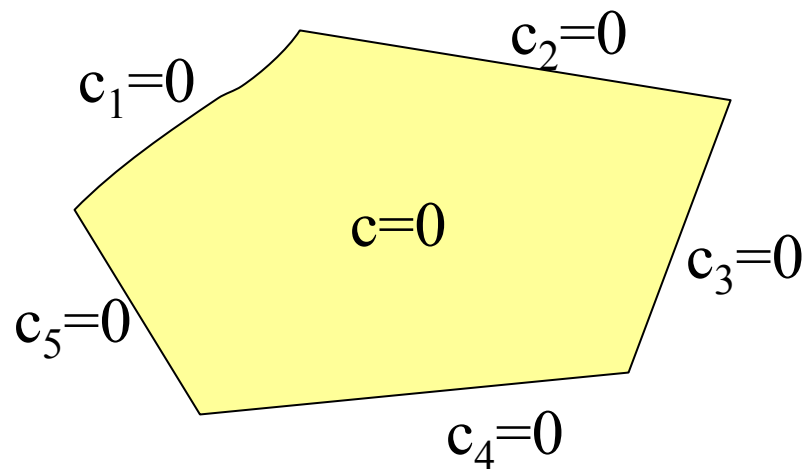
$$\vec{q} = -k\nabla T \quad (\text{Fourier's law for heat conduction})$$
$$\nabla \cdot \vec{q} = 0 \quad (\text{conservation law for heat}) \quad \longrightarrow \quad \nabla^2 T = 0$$

(steady heat flow)

Electrostatics

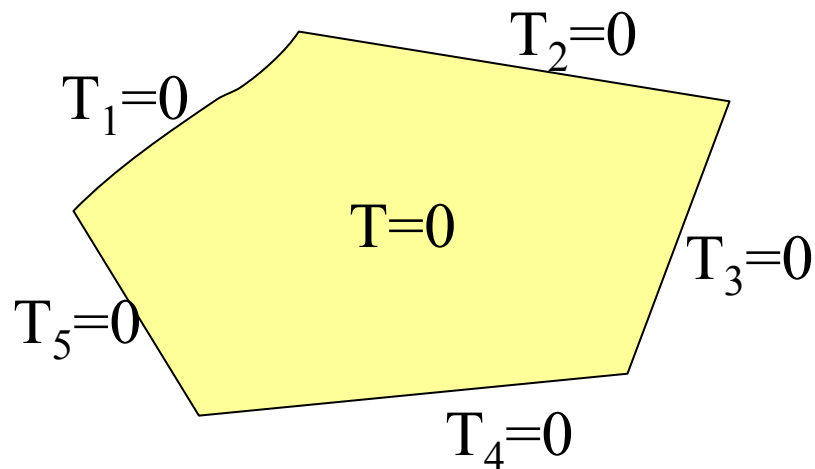


Steady state diffusion  $\nabla^2 c = 0$



Thermal conduction

$$\nabla^2 T = 0$$



# Uniqueness of Solution

Let's assume two different solutions,  $\Phi_a$  and  $\Phi_b$

$$\nabla^2 \cdot \Phi_a = -\frac{\rho_e}{\epsilon}; \quad \Phi_a = \Phi_i \quad \text{on } S_i$$

$$\nabla^2 \cdot \Phi_b = -\frac{\rho_e}{\epsilon}; \quad \Phi_b = \Phi_i \quad \text{on } S_i$$

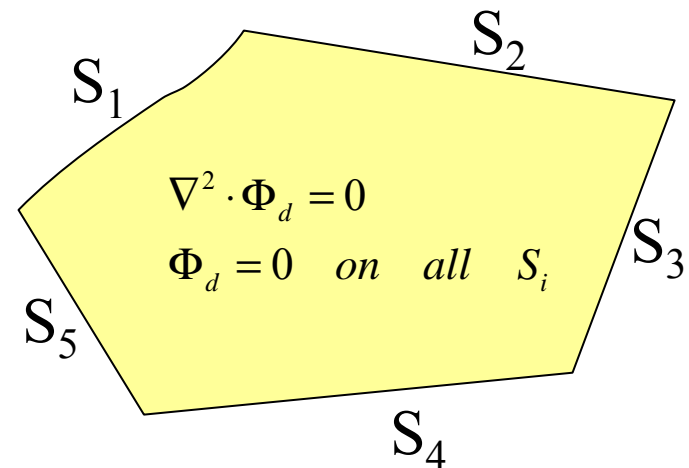
Then define  $\Phi_d = \Phi_a - \Phi_b$

$$\nabla^2 \cdot \Phi_d = 0; \quad \Phi_d = 0 \quad \text{on } S_i \quad (\text{satisfy Laplace Eq.})$$

Answer:

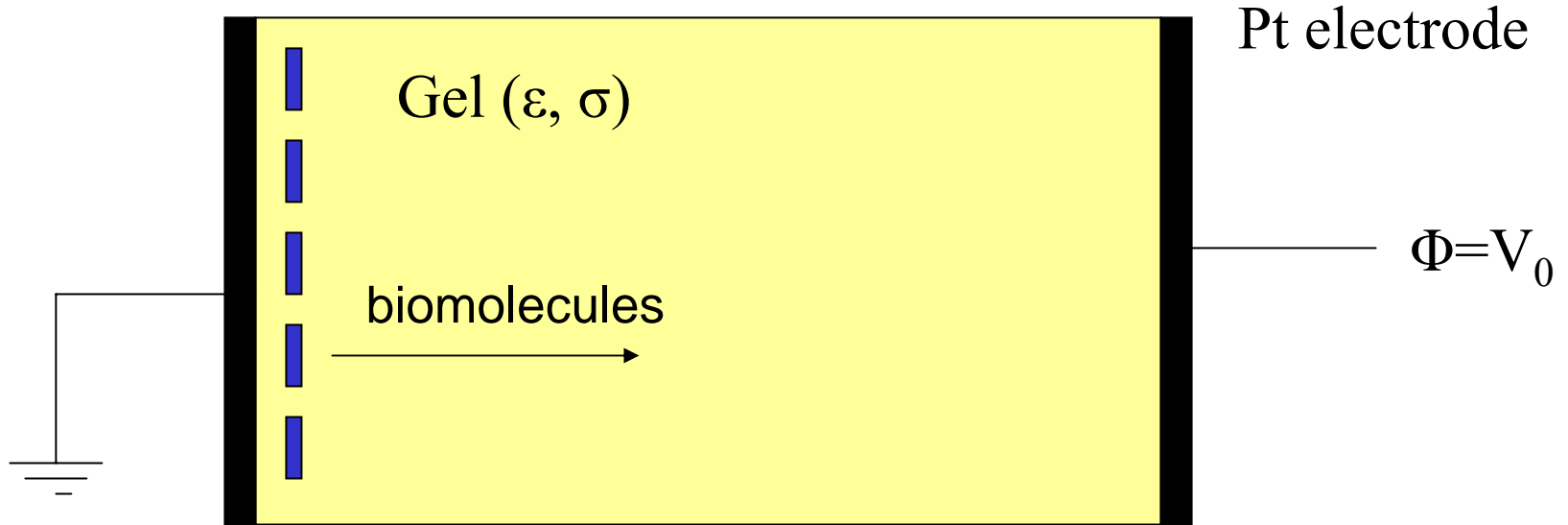
$$\Phi_d = 0 \quad \text{for everywhere}$$

$$\therefore \Phi_a - \Phi_b = 0$$



# Gel Electrophoresis

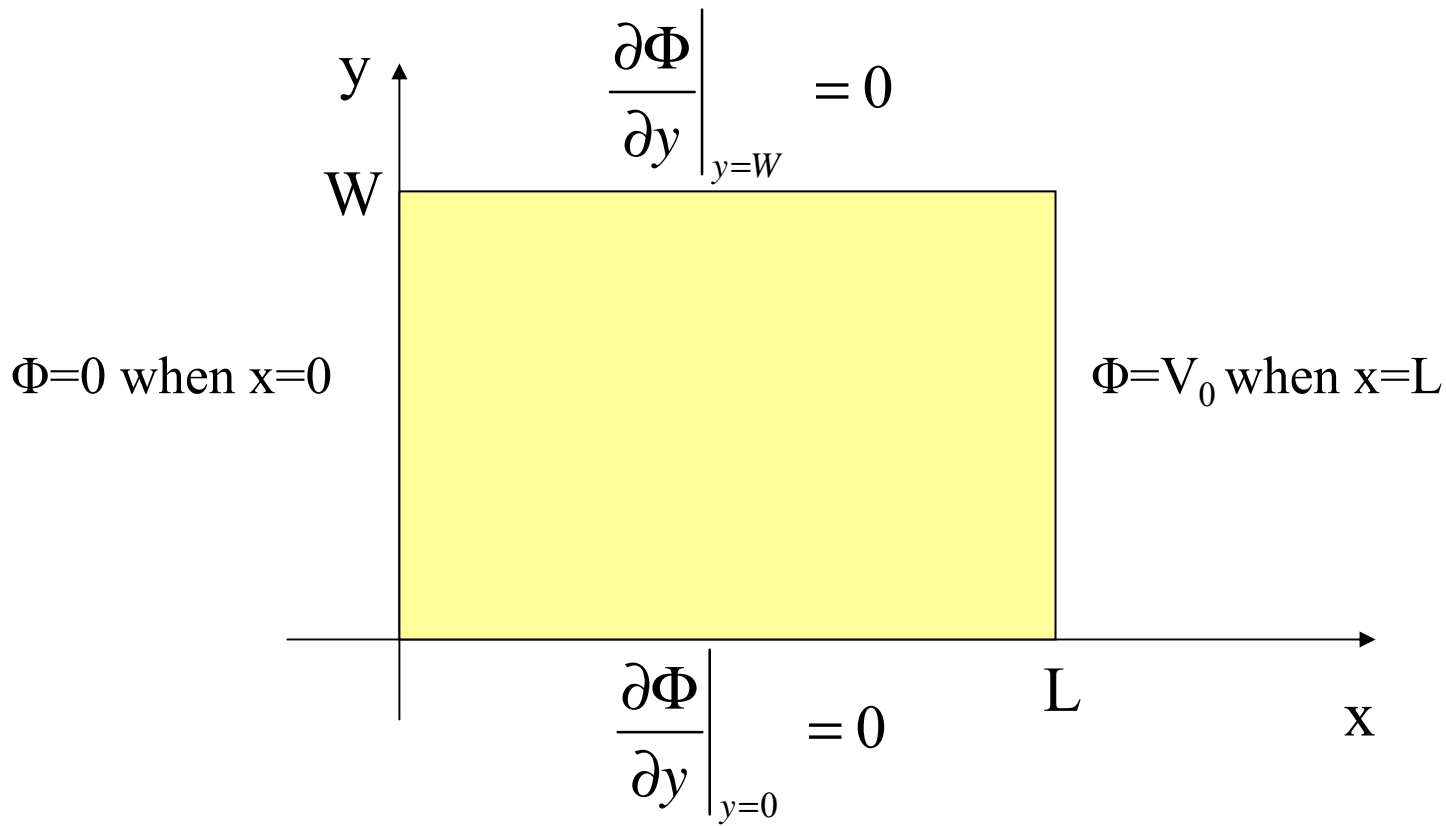
Plastic ( $\sigma = 0$ )



$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = 0 \text{ (electrostatics)} \quad \Rightarrow \quad \vec{E} = -\nabla \Phi$$

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t} = 0 \text{ (steady state, no charge accumulation)}$$

$$\nabla \cdot \vec{J} = \nabla \cdot (\sigma \vec{E}) = 0 \quad \Rightarrow \quad \nabla \cdot \vec{E} = 0 \quad \Rightarrow \quad \nabla^2 \Phi = 0$$



$\nabla \cdot \vec{J} = 0$   $\Rightarrow$  (no charge accumulation)

$\vec{J} = J_x \hat{x}$   
 $\longrightarrow$

$\Rightarrow J_y = \sigma E_y = 0 = \frac{\partial \Phi}{\partial y} \Big|_{y=0 \text{ or } W}$

$J=0$  (insulator)

# Boundary Conditions (For EQS approximation)

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$$



$$\hat{n} \cdot (\epsilon_1 \vec{E}_1 - \epsilon_2 \vec{E}_2) = \sigma_s$$

$$\nabla \times \vec{E} = 0$$

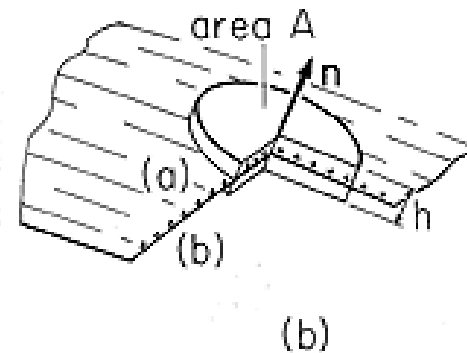
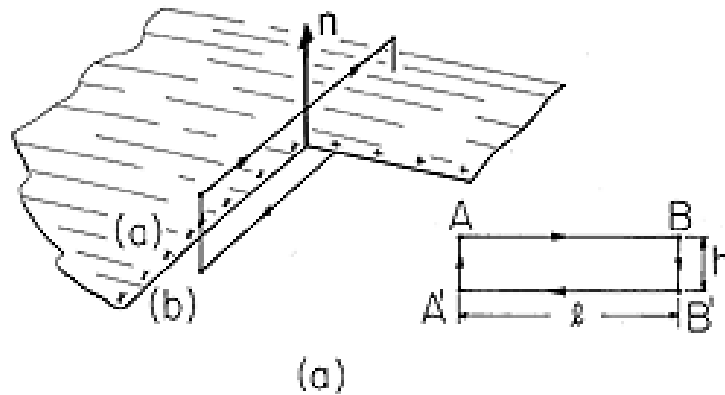


$$\hat{n} \times \vec{E}_1 = \hat{n} \times \vec{E}_2 \quad (\vec{E}_1|_{\text{tangential}} = \vec{E}_2|_{\text{tangential}})$$

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$



$$\hat{n} \cdot (\sigma_1 \vec{E}_1 - \sigma_2 \vec{E}_2) = -\frac{\partial \sigma_s}{\partial t}$$

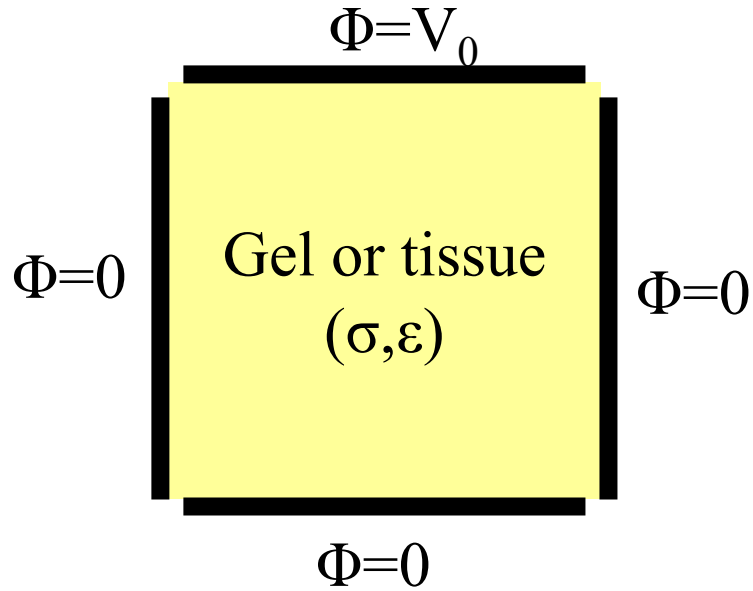


**Figure 5.3.1** (a) Differential contour intersecting surface supporting surface charge density. (b) Differential volume enclosing surface charge on surface having normal  $\mathbf{n}$ .

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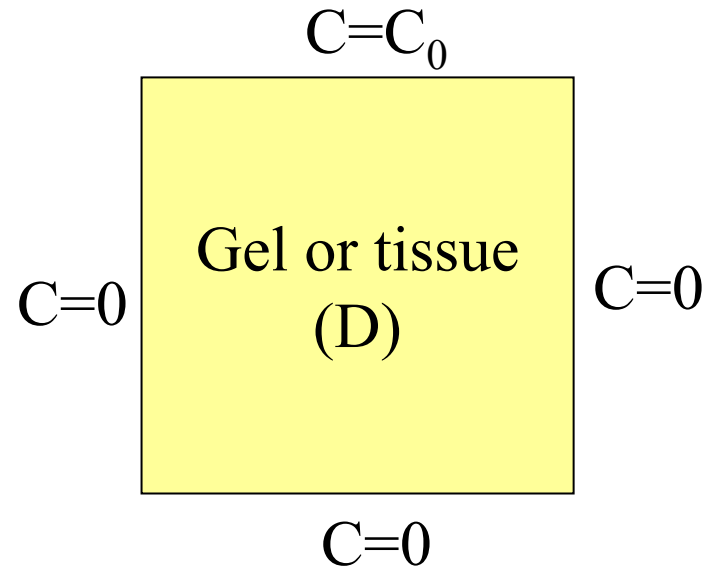
## Electrostatics



$$\nabla^2 \Phi = 0$$

$$\vec{J}_e = -\sigma \nabla \Phi$$

## Steady Diffusion

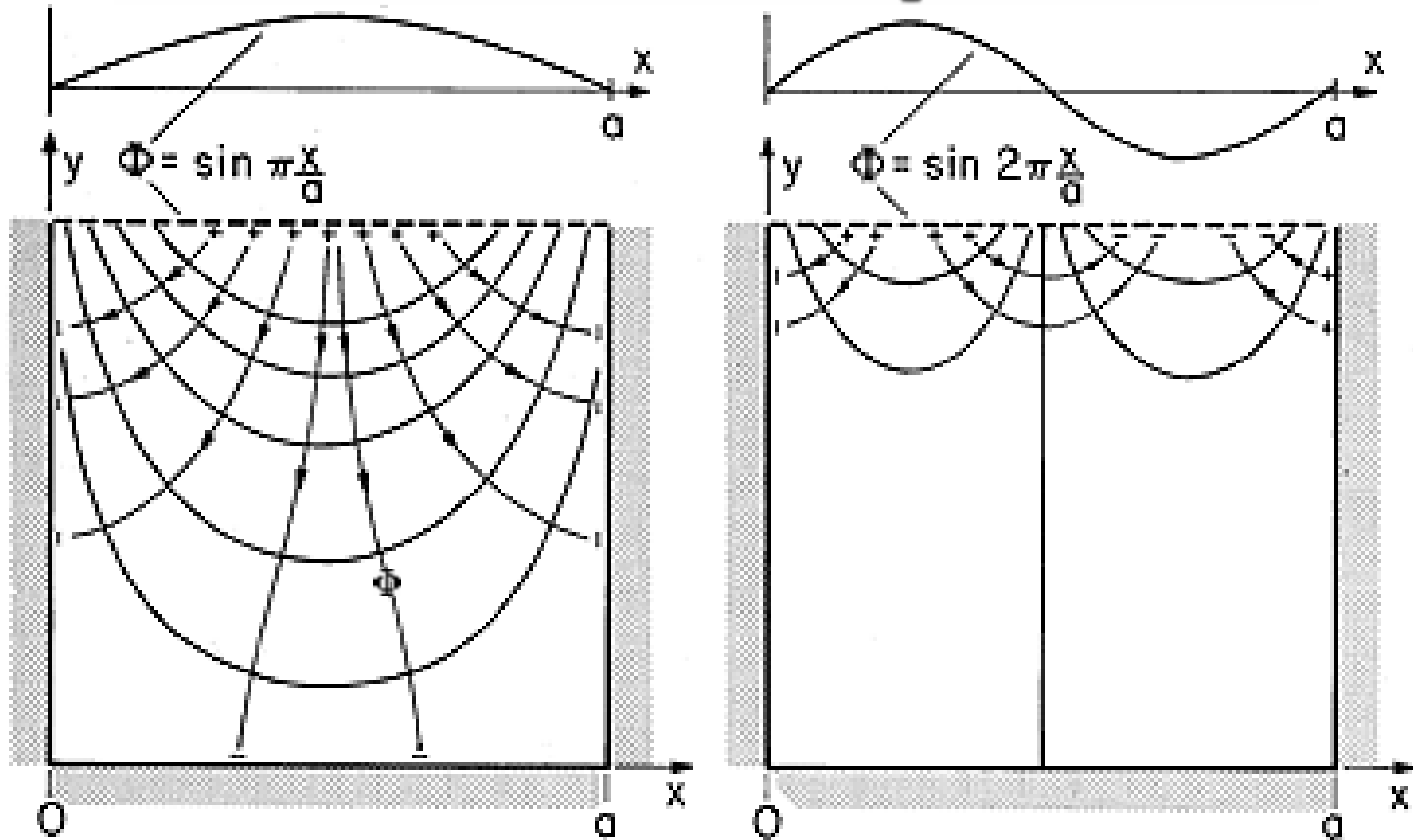


$$\nabla^2 C = 0$$

$$\vec{J}_i = -D_i \nabla c_i$$

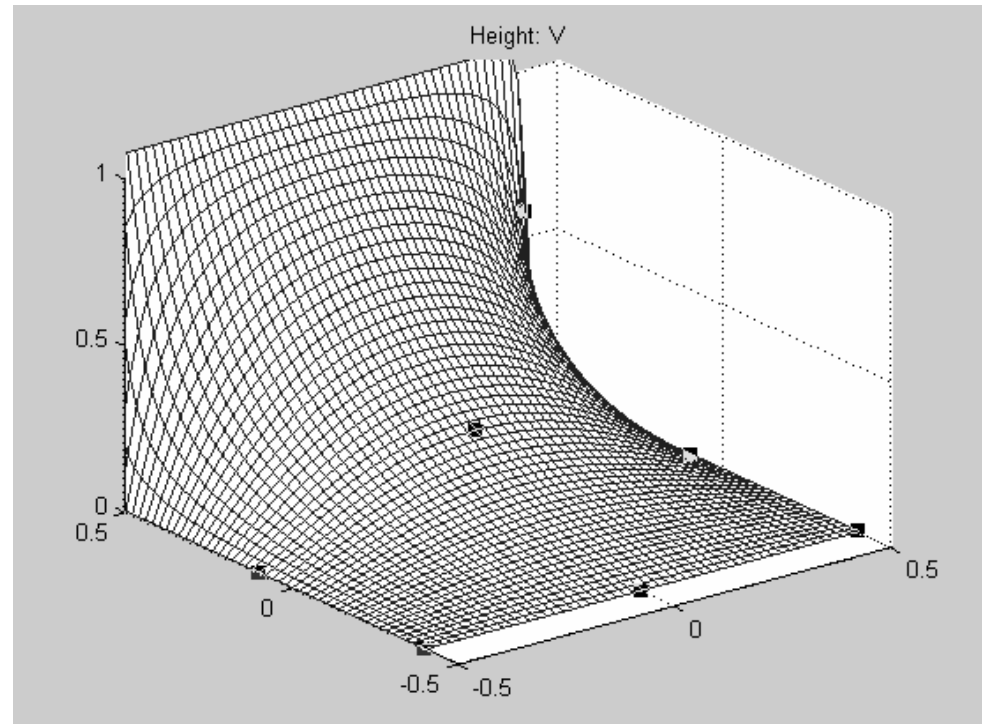
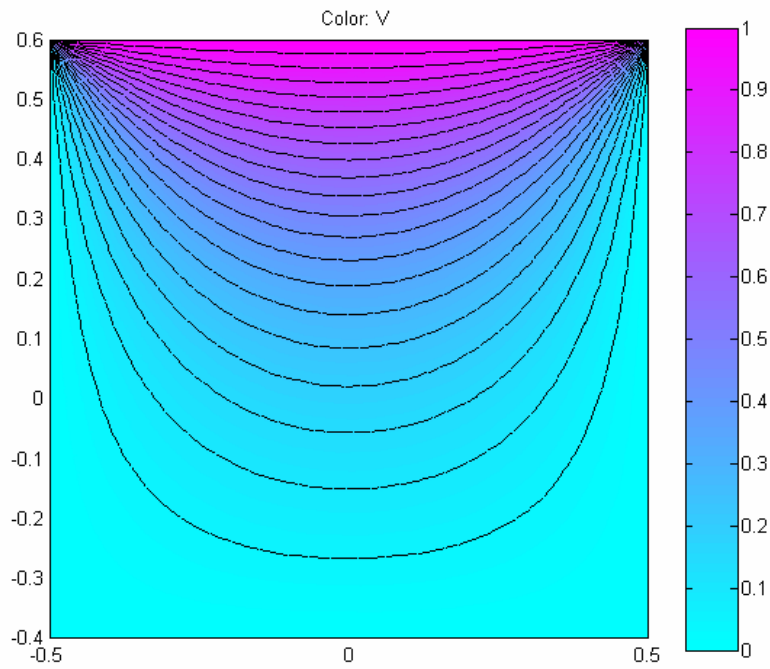


$$\Phi = A \sin kx \sinh ky$$



**Figure 5.5.1** Two of the infinite number of potential functions having the form of (1) that will fit the boundary conditions  $\Phi = 0$  at  $y = 0$  and at  $x = 0$  and  $x = a$ .

# Solution



# Known Solutions for Laplace equations

## Cylindrical Coordinates

$$\nabla^2 \Phi(\rho, \varphi, z) = 0 \Rightarrow \frac{\partial^2 \Phi}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial \Phi}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \varphi^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

$$\Phi(\rho, \varphi, z) = R(\rho)\Psi(\varphi)Z(z)$$

$$R(\rho) \Rightarrow \text{Bessel Functions } (J_n, N_n, I_n, K_n)$$

$$\Psi(\varphi) \Rightarrow \text{Trigonometric } (\sin, \cos, \sinh, \cosh)$$

$$Z(z) \Rightarrow \text{Trigonometric } (\sin, \cos, \sinh, \cosh)$$

## Spherical Coordinates

$$\nabla^2 \Phi(r, \theta, \varphi) = 0 \Rightarrow \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \Phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \varphi^2} = 0$$

$$\Phi(r, \theta, \varphi) = R(r)\Theta(\theta)\Psi(\varphi)$$

$$R(r) \Rightarrow \text{Spherical Bessel Functions}$$

$$\Theta(\theta) \Rightarrow \text{Legendre Functions } (P_n(\cos \theta))$$

$$\Psi(\varphi) \Rightarrow \text{Trigonometric } (\sin \varphi, \cos \varphi)$$

# Solving Laplace's Equation (Numerically)

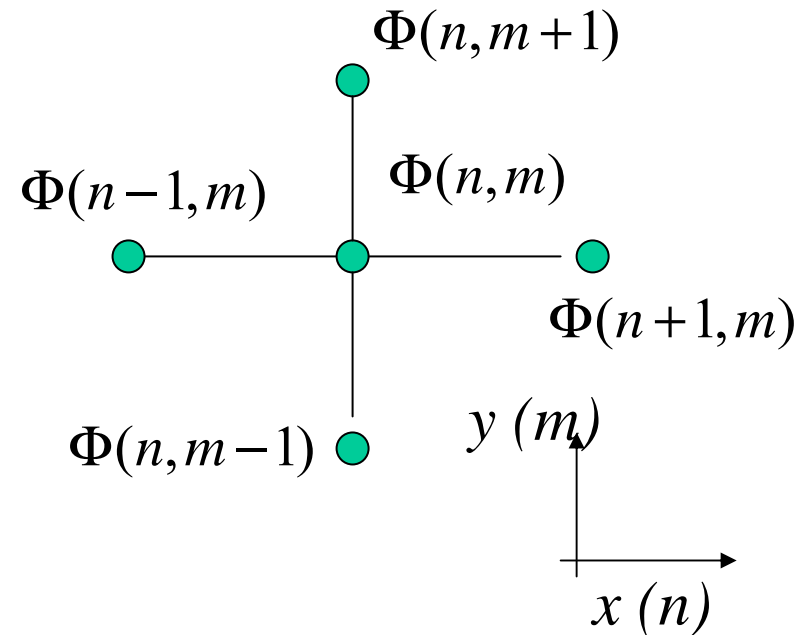
1D case:  $\frac{d^2\Phi}{dx^2} = 0 \rightarrow \Phi(x) = ax + b$

2D case:  $\frac{\partial^2\Phi}{\partial x^2} + \frac{\partial^2\Phi}{\partial y^2} = 0$

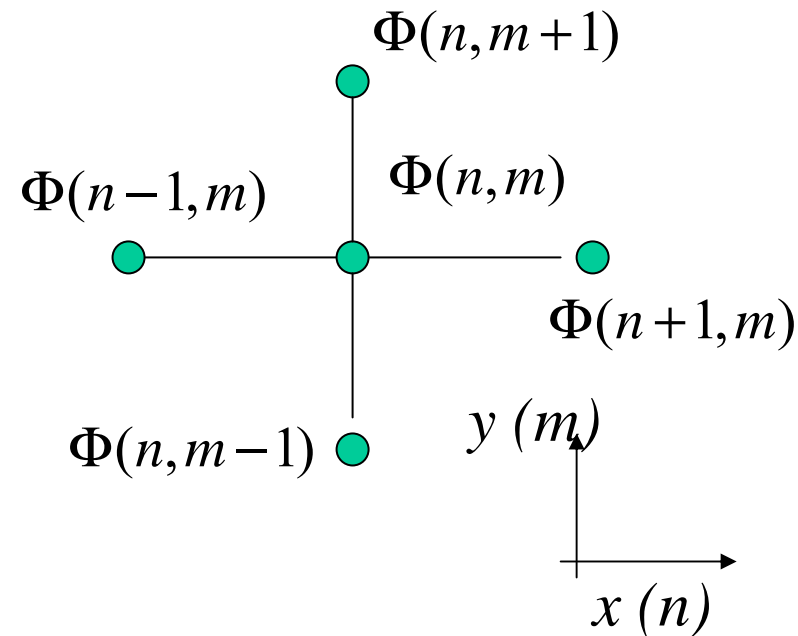
$$\frac{\partial\Phi}{\partial x}\left(n + \frac{1}{2}, m\right) = \Phi(n+1, m) - \Phi(n, m)$$

$$\frac{\partial\Phi}{\partial x}\left(n - \frac{1}{2}, m\right) = \Phi(n, m) - \Phi(n-1, m)$$

$$\frac{\partial^2\Phi}{\partial x^2}(n, m) = \frac{\partial\Phi}{\partial x}\left(n + \frac{1}{2}, m\right) - \frac{\partial\Phi}{\partial x}\left(n - \frac{1}{2}, m\right) = \Phi(n+1, m) + \Phi(n-1, m) - 2\Phi(n, m)$$



# Laplace's equation In discretized form



$$\frac{\partial^2 \Phi}{\partial x^2}(n, m) + \frac{\partial^2 \Phi}{\partial y^2}(n, m) =$$

$$\Phi(n+1, m) + \Phi(n-1, m) + \Phi(n, m+1) + \Phi(n, m-1) - 4\Phi(n, m) = 0$$

$$\Phi(n, m) = \frac{\Phi(n+1, m) + \Phi(n-1, m) + \Phi(n, m+1) + \Phi(n, m-1)}{4}$$

Value in the middle = average of surrounding values

# Finite Element Method

