

# NONLINEAR MECHANICAL SYSTEMS

## CANONICAL TRANSFORMATIONS AND NUMERICAL INTEGRATION

### Jacobi Canonical Transformations

**A Jacobi canonical transformation yields a Hamiltonian that depends on only one of the conjugate variable sets.**

**Assume dependence on new momentum alone.**

$$H(p^*, q^*) = K(p^*)$$

$$\partial K(p^*) / \partial q^* = 0$$

**Thus**

$$dp^* / dt = e^*$$

$$dq^* / dt = \partial K(p^*) / \partial p^* - f^*$$

**The simple relation between effort and the rate of change of momentum is recovered in the new coordinates.**

**EXAMPLE: SIMPLE HARMONIC OSCILLATOR**

**Hamiltonian**

$$H(p,q) = \frac{1}{2}(p^2/I + q^2/C)$$

**Hamilton's equations**

$$dq/dt = \partial H/\partial p = p/I$$

$$dp/dt = -\partial H/\partial q = -q/C$$

**Change variables from old (q,p) to new (P,Q)**

**Define  $Z_0 = \sqrt{IC}$  and the generating function**

$$S(q,Q) = Z_0(q^2/2) \cot Q$$

**The transformation equations are**

$$p = \partial S/\partial q = Z_0 q \cot Q$$

$$P = -\partial S/\partial Q = Z_0(q^2/2)/\sin^2 Q$$

**Express the old variables in terms of the new**

$$p = \sqrt{2P} \cos Q \sqrt{Z_0}$$

$$q = \sqrt{2P} \sin Q (1/\sqrt{Z_0})$$

**Define  $\omega_0 = \sqrt{1/IC}$  and the new Hamiltonian is**

$$H(P,Q) = \omega_0 P = K(P)$$

**Hamilton's equations in the new coordinates**

$$dQ/dt = \partial K/\partial P = \omega_0$$

$$dP/dt = -\partial K/\partial Q = 0$$

**Their solution is**

$$Q(t) = \omega_0 t + \text{constant}$$

$$P(t) = \text{constant}$$

**In essence this variable change has integrated the equations.**

**As the product of P and Q has the units of action (energy by time) it is sometimes called a (simple harmonic) actional transformation.**

**PHYSICAL INTERPRETATION:**

**P is proportional to the total system energy.**

**Its square root is proportional to oscillation amplitude.**

**Q is the phase angle of the oscillations.**

**In general, finding Jacobi canonical transformations requires solving a non-trivial partial differential equation.**

**A practical alternative is to separate the Hamiltonian into two parts, one with a known Jacobi canonical transform.**

$$H(p,q) = H_j(p,q) + H_n(p,q)$$

**Apply the known Jacobi canonical transformation**

$$H^*(P,Q) = H_j^*(P) + H_n^*(P,Q)$$

**We may represent the second term as a set of canonical forces**

$$e^*(P,Q) = -\partial H_n^*/\partial Q$$

$$f^*(P,Q) = -\partial H_n^*/\partial P$$

**The transformed equations become**

$$dP/dt = e^*(P,Q)$$

$$dQ/dt = \partial H_j^*/\partial P - f^*(P,Q)$$

**An advantage of this change of variables is that, in effect, it integrates the fundamental oscillatory mode of the solution.**

**EXAMPLE: SIMPLE PENDULUM**

**For large amplitudes, the simple pendulum is a nonlinear oscillator.**

$$H(\eta, \theta) = \eta^2/2 + 1 - \cos \theta$$

where

$\theta$  angle with respect to the vertical

$\eta$  corresponding angular momentum

**Expand the cosine as a power series**

$$H(\eta, \theta) = \eta^2/2 + \theta^2/2 - \theta^4/4! + \theta^6/6! - \dots$$

**The Hamiltonian is quadratic in momentum and displacement with additional terms in displacement of fourth power and higher.**

**Until the fourth power of the angle in radians becomes significant,**

**the nonlinear pendulum may be treated as linear system**

**with a Hamiltonian that is quadratic in momentum and displacement.**

**For the quadratic terms have a known Jacobi canonical transformation: the simple harmonic actional.**

**Split the Hamiltonian as follows**

$$H(\eta, \theta) = \eta^2/2 + \theta^2/2 + (1 - \cos\theta - \theta^2/2)$$

$$H(\eta, \theta) = K(\eta, \theta) + N(\eta, \theta)$$



**Apply the simple harmonic actional**

$$\theta = \sqrt{2P} \sin Q$$

$$\eta = \sqrt{2P} \cos Q$$

**The Hamiltonian becomes**

$$H^*(P, Q) = K^*(P) + N^*(P, Q)$$

**In the original variables, the system equations are**

$$d\eta/dt = -\partial H/\partial\theta$$

$$d\theta/dt = \partial H/\partial\eta$$

**In the new variables, the system equations become**

$$dP/dt = -\partial N^*/\partial Q$$

$$dQ/dt = 1 + \partial N^*/\partial P$$

**Transformation does not change the value of either K or N.**

**Use the chain rule on the original N which depends only on  $\theta$ .**

$$\partial N^*/\partial Q = (\partial N/\partial \theta) (\partial \theta/\partial Q)$$

$$\partial N^*/\partial P = (\partial N/\partial \theta) (\partial \theta/\partial P)$$

$$\partial N/\partial \theta = \sin \theta - \theta$$

$$\partial \theta/\partial Q = \sqrt{2P} \cos Q$$

$$\partial \theta/\partial P = \sin Q (1/\sqrt{2P} )$$

**The transformed equations become**

$$dP/dt = [\sqrt{2P} \sin Q - \sin(\sqrt{2P} \sin Q)][\sqrt{2P} \cos Q]$$

$$dQ/dt = 1 + [\sin(\sqrt{2P} \sin Q) - \sqrt{2P} \sin Q][\sin Q (1/\sqrt{2P} )]$$

**Use the transformation equations to express the rates of change as a function of both old and new variables.**

$$dP/dt = (\theta - \sin\theta)\eta$$

$$dQ/dt = 1 + (\sin\theta - \theta)\theta/2P$$

$$\theta = \sqrt{2P} \sin Q$$

$$\eta = \sqrt{2P} \cos Q$$

**What have we gained?**

**The system equations are simpler in the old variables**

$$d\eta/dt = -\sin\theta$$

$$d\theta/dt = \eta$$

**In the new variables, the solution is far more stable numerically.**

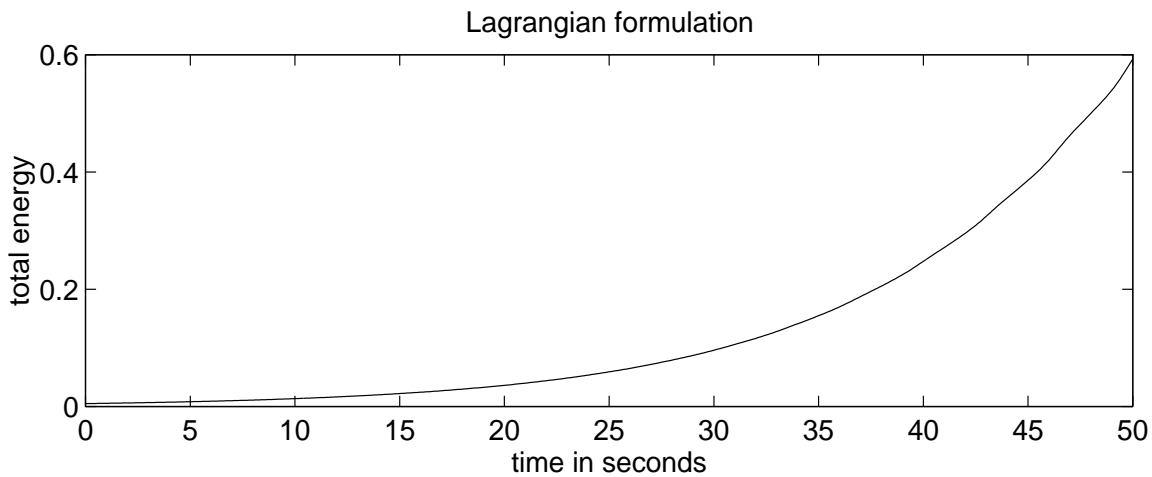
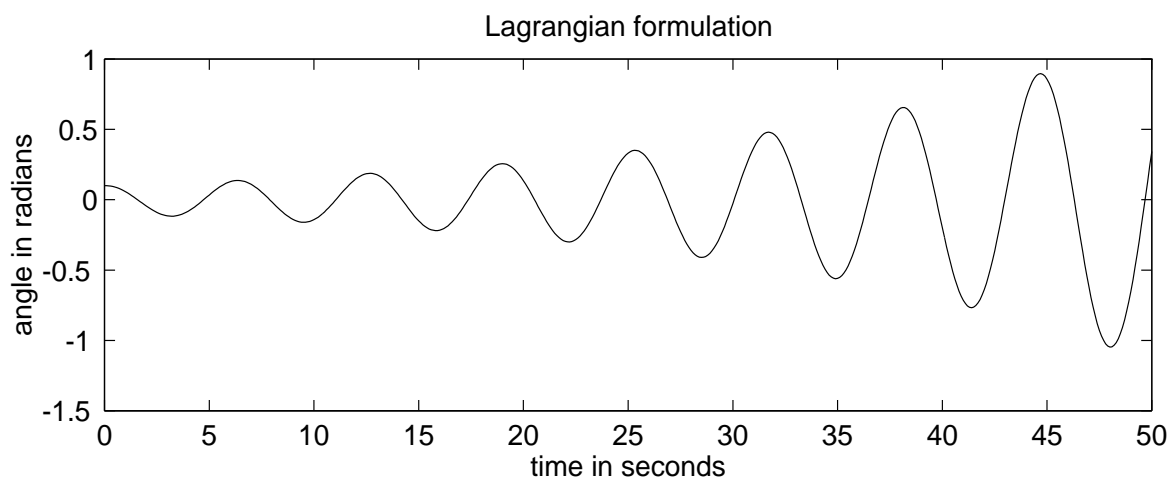
## Simple Euler integration algorithm

starting time 0 seconds, final time 50 seconds, time step 0.1 seconds.

start from rest at an angle of 0.1 radians ( $\approx 6^\circ$ )

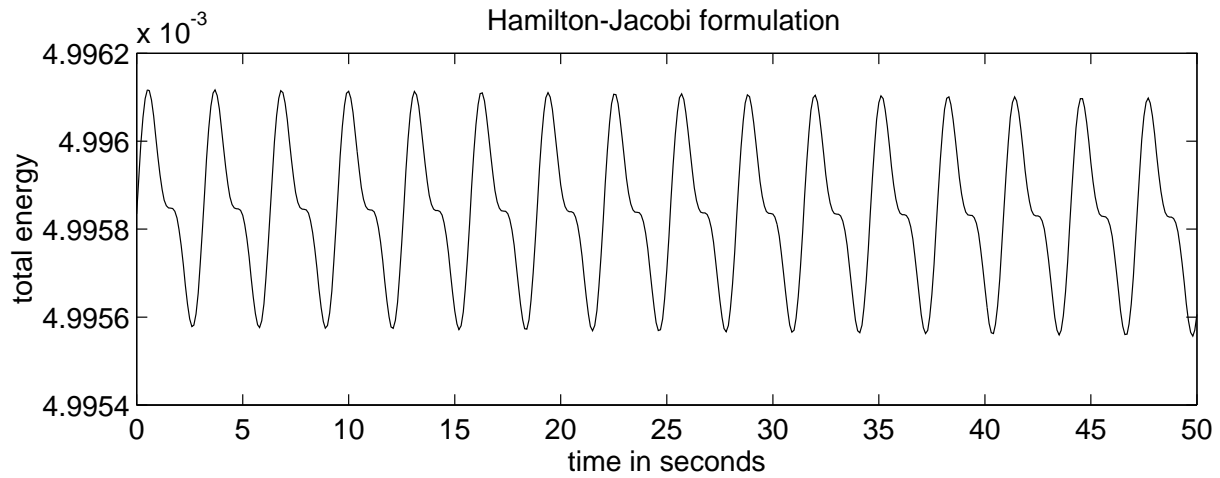
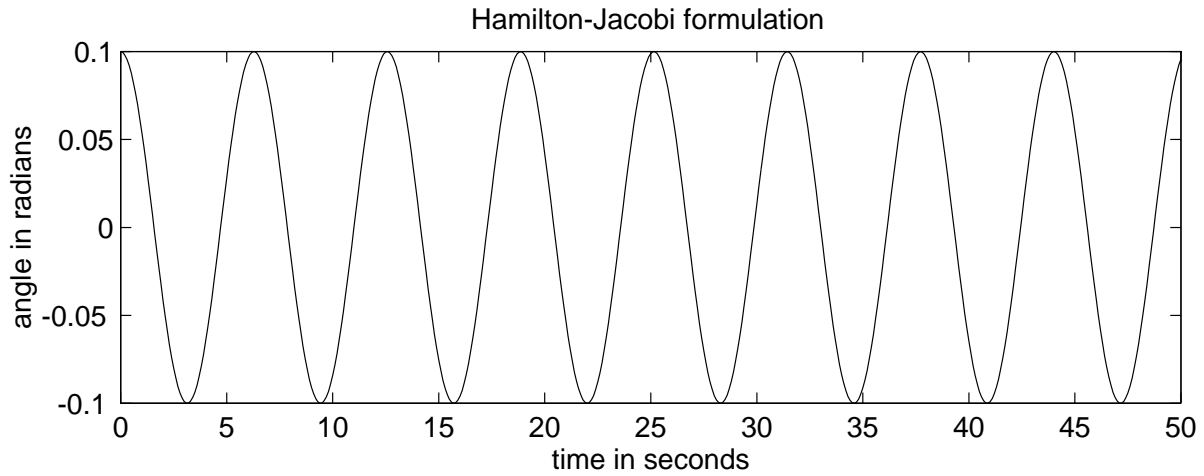
In old coordinates, simulation is unstable.

Total system energy grows exponentially.

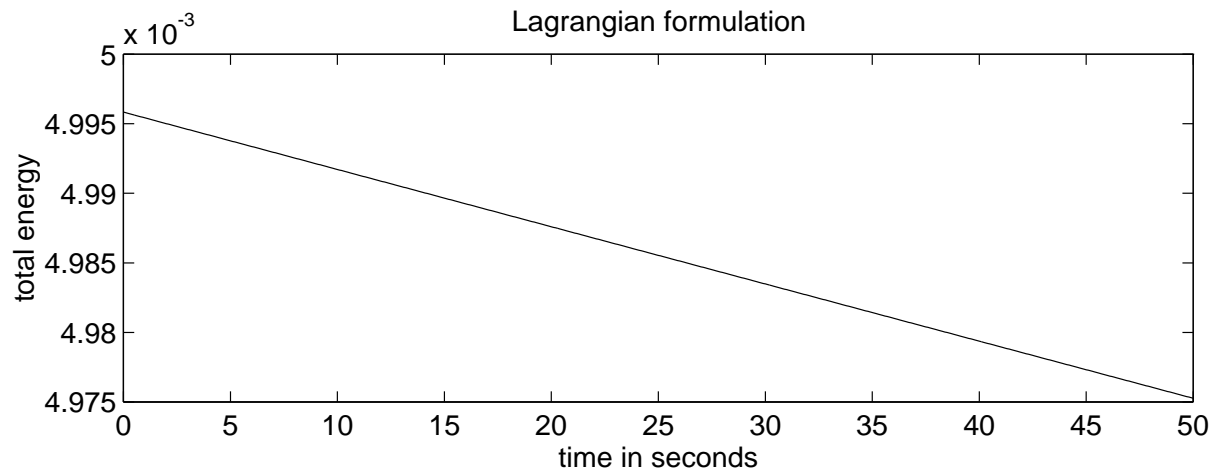
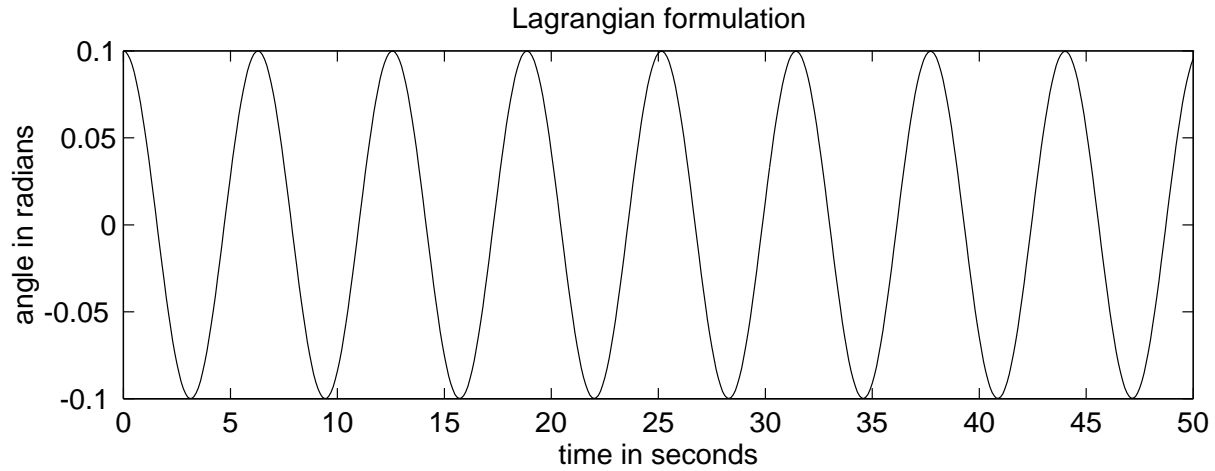


**In new coordinates, the simulation is stable.**

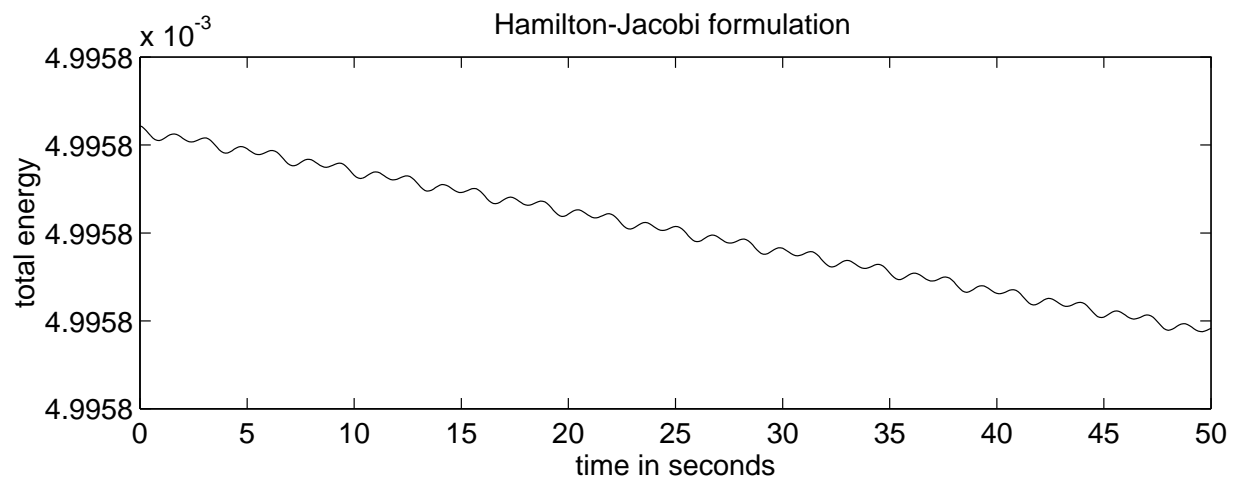
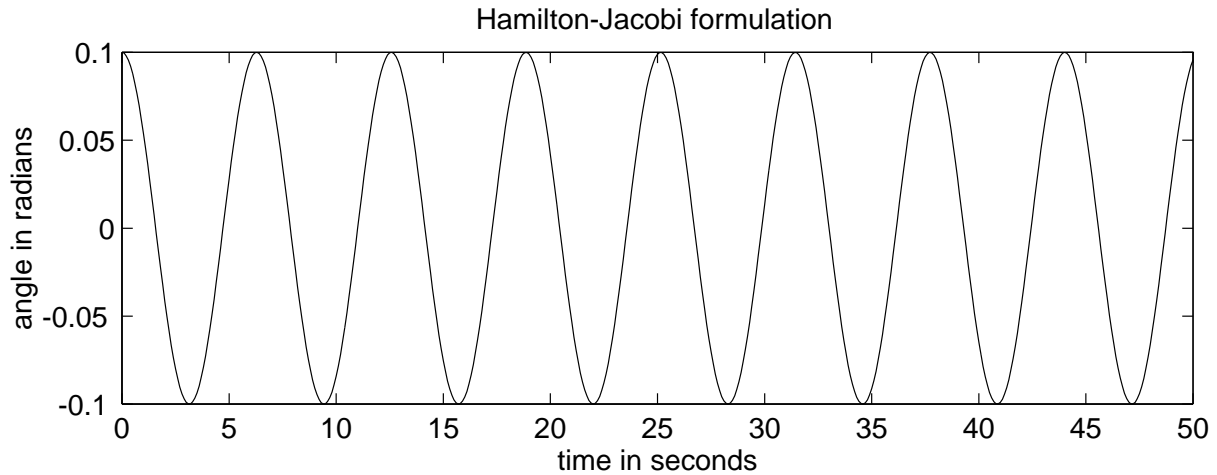
**Total system energy variation:  $5.6 \times 10^{-7}$ .**



**Perform the same integration using a third-order fixed-step Runge-Kutta algorithm.**

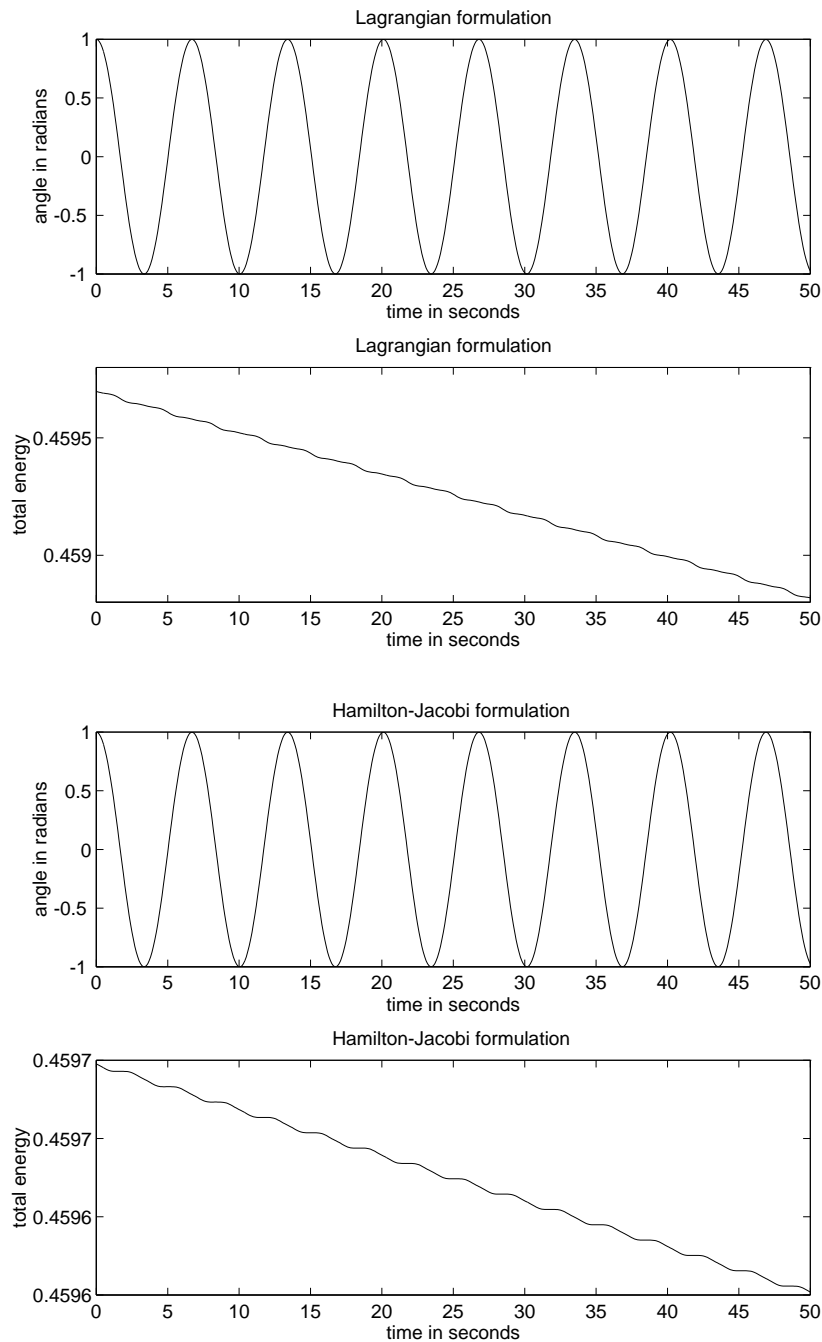


**Total energy, declines steadily by  $2.1 \times 10^{-5}$  over 50 seconds.**



**Total energy also decreases, but by  $2.3 \times 10^{-10}$  — a hundred thousand time less.**

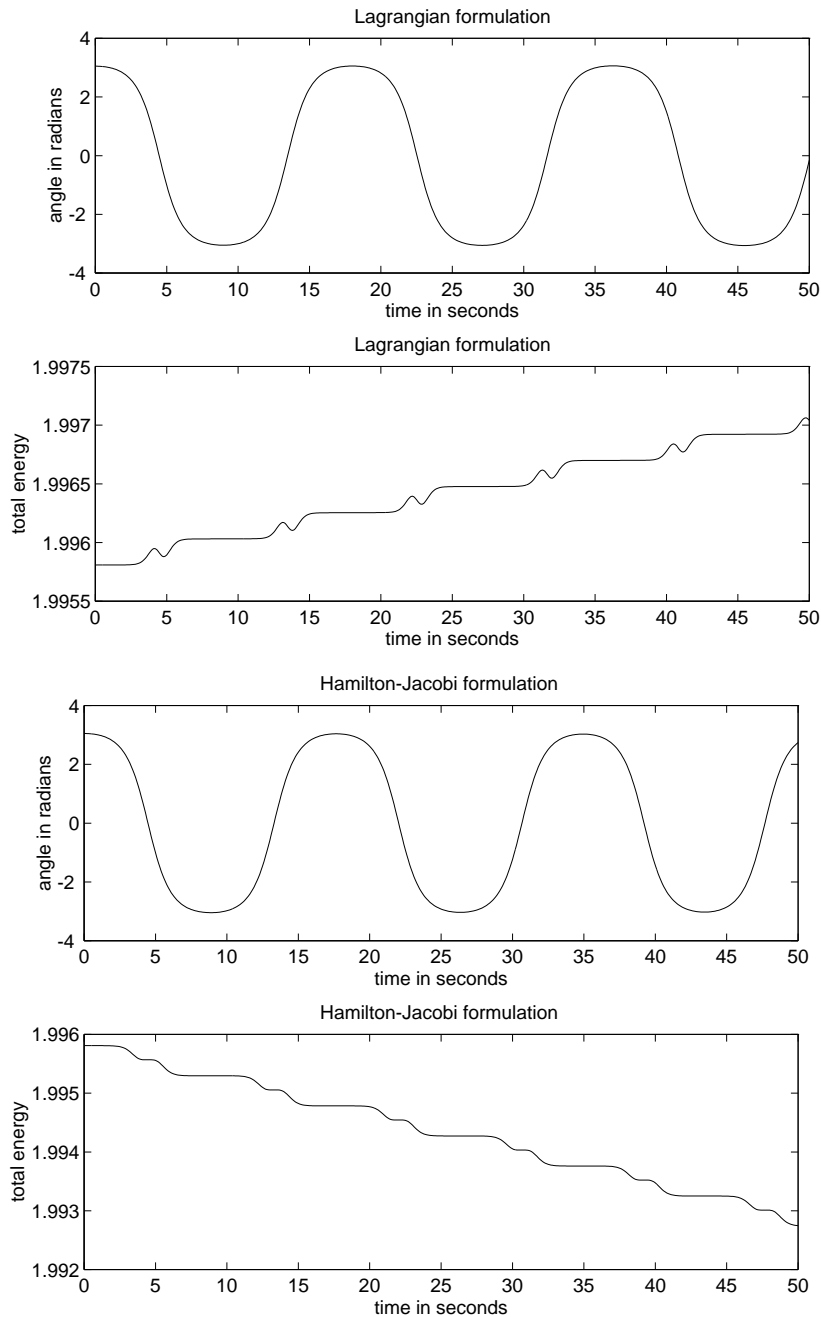
Start the pendulum from rest at 1 radian ( $\approx 57^\circ$ ) and use the same integration algorithm and parameters



Again, the transformed equations produce a smaller decline in energy, though the difference is less pronounced –  $8.8 \times 10^{-4}$  vs.  $1.5 \times 10^{-4}$ .

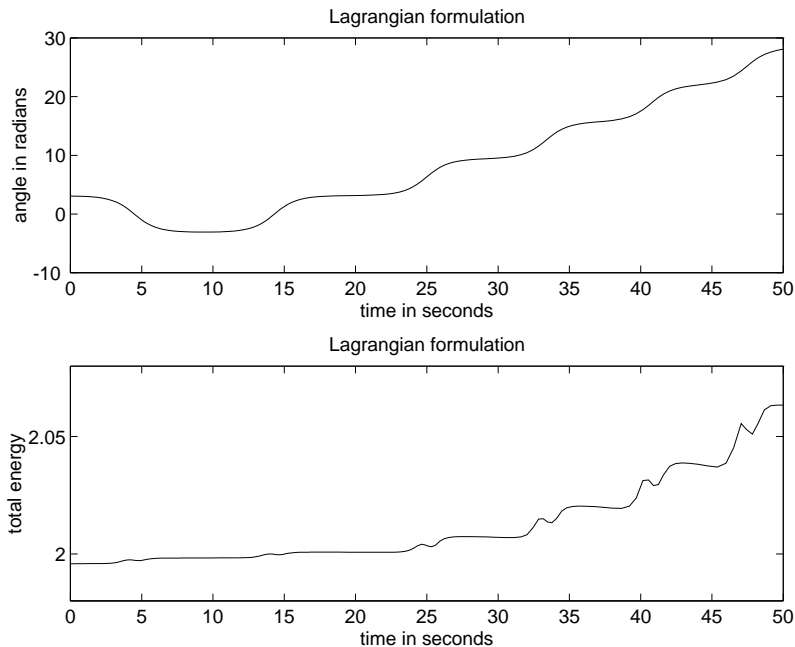


Start from rest at  $5^\circ$  off vertically upright (3.05 radians) and use the same integration algorithm and parameters



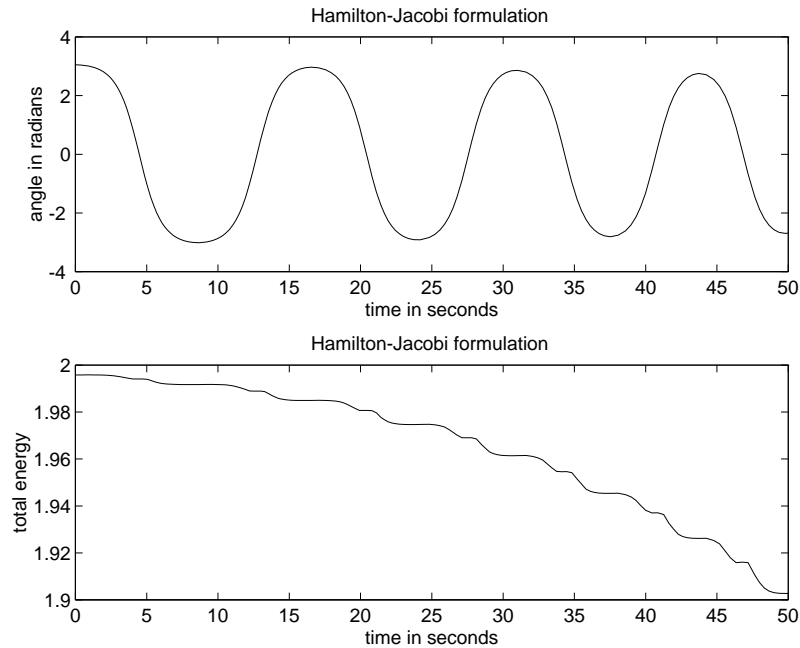
Now the original formulation is unstable – energy increases by  $1.3 \times 10^{-3}$  in 50 seconds. The transformed equations yield a decline of energy of  $3.1 \times 10^{-3}$ .

Start from the same initial conditions but use MATLAB's ode23, a 3rd-order adaptive Runge-Kutta algorithm with error tolerance of  $1.0 \times 10^{-3}$



Now the steady increase in computed total energy in the original formulation results in a major departure of the computed angle from what it should be

– the simulation claims that after one oscillation the pendulum will spin continuously in one direction.



The transformed equations do not exhibit this behavior, though the computed energy declines substantially ( $9.3 \times 10^{-2}$  in 50 seconds).

## **POINTS:**

- **Never believe anything you get from a computer. Find some way of cross checking the results. One effective method is to compute a known invariant, in this case energy.**
- **The equations in the original variables may look simpler, but that is deceptive. In fact the transformed equations have been partially integrated by the transformation and so present a less demanding task to the integration algorithm.**
- **A little analysis up front can have a dramatic effect on the accuracy of numerical computations.**