

3 Conservation of volume in phase space

We show (via the example of the pendulum) that frictionless systems *conserve* volumes (or areas) in phase space.

Conversely, we shall see, dissipative systems *contract* volumes.

Suppose we have a 3-D phase space, such that

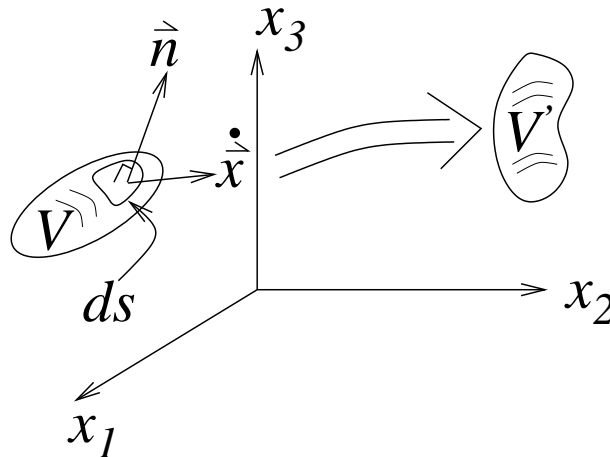
$$\begin{aligned}\dot{x}_1 &= f_1(x_1, x_2, x_3) \\ \dot{x}_2 &= f_2(x_1, x_2, x_3) \\ \dot{x}_3 &= f_3(x_1, x_2, x_3)\end{aligned}$$

or

$$\frac{d\vec{x}}{dt} = \vec{f}(\vec{x})$$

The equations describe a “flow,” where $d\vec{x}/dt$ is the velocity.

A set of initial conditions enclosed in a volume V flows to another position in phase space, where it occupies a volume V' , neither necessarily the same shape nor size:



Assume the volume V has surface S .

Let

- ρ = density of initial conditions in V ;
- $\rho \vec{f}$ = rate of flow of points (trajectories emanating from initial conditions) through unit area perpendicular to the direction of flow;
- ds = a small region of S ; and
- \vec{n} = the unit normal (outward) to ds .

Then

$$\text{net flux of points out of } S = \int_S (\rho \vec{f} \cdot \vec{n}) ds$$

or

$$\int_V \frac{\partial \rho}{\partial t} dV = - \int_S (\rho \vec{f} \cdot \vec{n}) ds$$

i.e., a positive flux \implies a loss of “mass.”

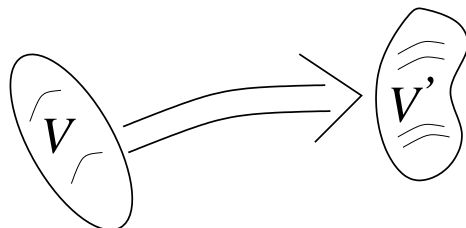
Now we apply the divergence theorem to convert the integral of the vector field $\rho \vec{f}$ on the surface S to a volume integral:

$$\int_V \frac{\partial \rho}{\partial t} dV = - \int_V [\vec{\nabla} \cdot (\rho \vec{f})] dV$$

Letting the volume V shrink, we have

$$\frac{\partial \rho}{\partial t} = -\vec{\nabla} \cdot (\rho \vec{f})$$

Now follow the motion of V to V' in time δt :



The boundary deforms, but it always contains the same points.

We wish to calculate $d\rho/dt$, which is the rate of change of ρ as the volume moves:

$$\begin{aligned} \frac{d\rho}{dt} &= \frac{\partial\rho}{\partial t} + \frac{\partial\rho}{\partial x_1} \frac{dx_1}{dt} + \frac{\partial\rho}{\partial x_2} \frac{dx_2}{dt} + \frac{\partial\rho}{\partial x_3} \frac{dx_3}{dt} \\ &= -\vec{\nabla} \cdot (\rho\vec{f}) + (\vec{\nabla}\rho) \cdot \vec{f} \\ &= -(\vec{\nabla}\rho) \cdot \vec{f} - \rho\vec{\nabla} \cdot \vec{f} + (\vec{\nabla}\rho) \cdot \vec{f} \\ &= -\rho\vec{\nabla} \cdot \vec{f} \end{aligned}$$

Note that the number of points in V is

$$N = \rho V$$

Since points are neither created nor destroyed we must have

$$\frac{dN}{dt} = V \frac{d\rho}{dt} + \rho \frac{dV}{dt} = 0.$$

Thus, by our previous result,

$$-\rho V \vec{\nabla} \cdot \vec{f} = -\rho \frac{dV}{dt}$$

or

$$\frac{1}{V} \frac{dV}{dt} = \vec{\nabla} \cdot \vec{f}$$

This is called the *Lie derivative*.

We shall next arrive at the following main results by example:

- $\vec{\nabla} \cdot \vec{f} = 0 \Rightarrow$ volumes in phase space are conserved. Characteristic of conservative or Hamiltonian systems.
- $\vec{\nabla} \cdot \vec{f} < 0 \Rightarrow dV/dt < 0 \Rightarrow$ volumes in phase space contract. Characteristic of dissipative systems.

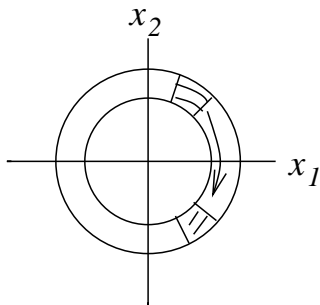
We use the example of the pendulum:

$$\begin{aligned} \dot{x}_1 &= \dot{x}_2 &= f_1(x_1, x_2) \\ \dot{x}_2 &= -\frac{g}{l} \sin x_1 &= f_2(x_1, x_2) \end{aligned}$$

Calculate

$$\vec{\nabla} \cdot \vec{f} = \frac{\partial \dot{x}_1}{\partial x_1} + \frac{\partial \dot{x}_2}{\partial x_2} = 0 + 0$$

Pictorially



Note that the area is conserved.

Conservation of areas holds for *all* conserved systems. This is conventionally derived from Hamiltonian mechanics and the canonical form of equations of motion.

In conservative systems, the conservation of volumes in phase space is known as *Liouville's theorem*.

4 Damped oscillators and dissipative systems

4.1 General remarks

We have seen how conservative systems behave in phase space. What about dissipative systems?

What is a fundamental difference between dissipative systems and conservative systems, aside from volume contraction and energy dissipation?

- Conservative systems are invariant under time reversal.
- Dissipative systems are not; they are *irreversible*.